

Electric flux density :

The flux due to the electric field 'E' can be calculated by using the definition of the flux.

According to eqn $E = \frac{Q}{4\pi\epsilon_0 R^2} aR$

the electric field intensity is dependent on the medium in which the charge is place.

New vector D is defined by

$$D = \epsilon_0 E \quad \dots\dots\dots (1)$$

The electric flux Ψ in terms of D ,

$$\Psi = \int_S D \cdot ds \quad \dots\dots\dots (2)$$

$D \rightarrow$ electric flux density $\rightarrow C/m^2$

$x \rightarrow$ measured in coulomb's

Electric flux density is also called electric displacement

Example : For an infinite sheet of charge

$$E = \frac{\rho_s}{2\epsilon_0} a_n \quad \dots\dots\dots (3)$$

$$D = \epsilon_0 E = \frac{\epsilon_0 \rho_s}{2\epsilon_0} a_n$$

$$D = \frac{\rho_s}{2} a_n \quad \dots\dots\dots (4)$$

For a volume charge distribution

$$a = \int_v \rho_v dv$$

$$D = \int_v \frac{\rho_v av}{4\pi R^2} aR \quad \dots\dots\dots (5)$$

From eqn (4) & (5) 'D' is a function of charge & position only it is independent of medium

Gauss's law – Maxwell's equation

Gauss's law states that the total electric flux Ψ through any closed surface is equal to the total charge enclosed by that surface.

$$\Psi = Q_{\text{enc}} \quad \dots\dots\dots (1)$$

$$\text{ie } \Psi = \oint d\Psi = \oint D \cdot ds$$

$$\text{total charge enclosed } Q = \int_v \rho_v dv$$

$$Q = \oint D \cdot ds = \int_v \rho_v dv \quad \dots\dots\dots (2)$$

By applying divergence theorem

$$\oint D \cdot ds = \int_v \nabla \cdot D dv \quad \dots\dots\dots (3)$$

Compare the two volume integrals in eqn (2) & (3)

$$\rho_v = \nabla \cdot D \quad \dots\dots\dots (4)$$

The above equation is a Maxwell first equation. The equation states that volume charge density is same as the divergence of electric flux density

ρ_v at any point – charge per unit volume at that point.

Note :

- Equation (A) & (4) state Gauss's law in different ways eqn (A) is the integral form
Eqn (4) is the differential (or) point form of Gauss's law
- Gauss's law is an alternative statement of Coulomb's law. Proper application of divergence theorem to Coulomb's law results Gauss's Law.
- By using Gauss's law E (or) D for symmetrical charge distributions such as point charge an infinite line charge an infinite cylindrical surface charge and a spherical distribution of charge.
- Rectangular symmetry depends on x (or y or z)
- Cylindrical symmetry depends on ρ
- Spherical symmetry depends on r (independent of θ & φ)

Gauss's law cannot be used to determine E (or) D. When the charge distribution is not symmetric.

Applications of Gauss's law :

- Gauss's law is used to calculate the electric field whether symmetry exists
- Surface of D is chosen such that D is normal or tangential to the Gaussian surface.

D is normal to the surface

$$D \cdot ds = D ds$$

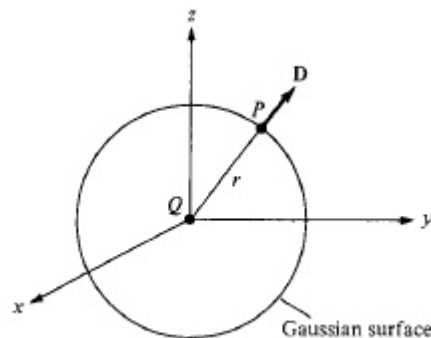
D is tangential to the surface

$$D \cdot ds = 0$$

Point charge :

A point charge Q is located at the origin. To determine D at point p. Here spherical surface containing 'p' satisfy the symmetry conditions is used.

Spherical surface centered at the origin is the Gaussian surface



Gaussian Surface

Fig25. Gaussian surface about a point charge

D is normal to the Gaussian surface

$$D = D_r \cdot a_r \quad \dots\dots\dots (1)$$

Gauss's law $Q = \oint_S D \cdot ds \quad \dots\dots\dots (2)$

$$Q = \oint_S D_r \cdot ds = D_r 4\pi r^2 \quad \dots\dots\dots (3)$$

$$\oint_S ds = 4\pi r^2 = \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin\theta d\theta d\varphi$$

$$D_r = D = \frac{Q}{4\pi r^2} \mathbf{a}_r \quad \dots\dots\dots (4)$$

Infinite line charge :

Infinite line of uniform charge ρ_L C/m along z- axis. To determine D at point p

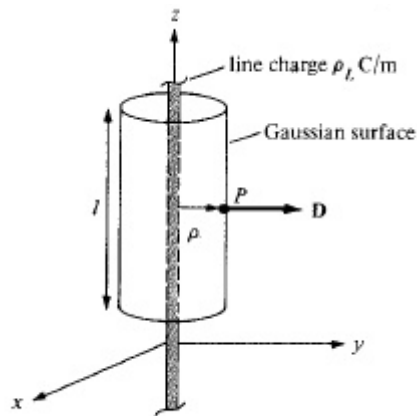


Fig26. Gaussian surface about an infinite line charge

The electric flux density 'D' is constant on and normal to the cylindrical Gaussian surface

$$D = D_\rho \mathbf{a}_\rho$$

Apply Gauss's law to an arbitrary length l of the line.

$$\begin{aligned} \rho_L L = Q &= \oint_S D \cdot d\mathbf{s} = D_\rho \oint_S d\mathbf{s} \\ &= D_\rho 2\pi\rho l \end{aligned}$$

$$\oint_S d\mathbf{s} = 2\pi\rho l \rightarrow \text{Surface area of the Gaussian surface.}$$

In the top & bottom of the surface 'D' is tangential to the Gaussian surface.

$$\oint_S D \cdot d\mathbf{s} = 0, \quad D_\rho = D = \frac{Q \rho_L l}{2\pi\rho l} = \frac{\rho_L}{2\pi\rho} \mathbf{a}_\rho$$

Infinite sheet of charge :

Infinite sheet of uniform charge $\rho_s \text{ C/m}^2$ lying on the $z=0$ plane. To determine 'D' at point. 'P'

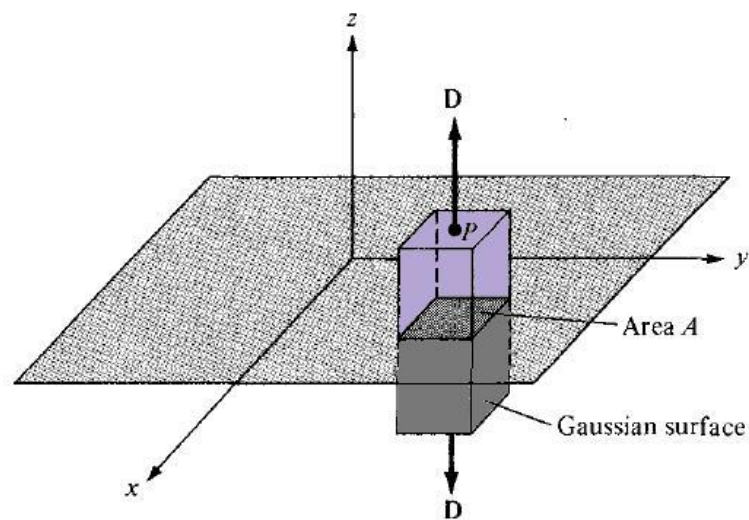


Fig6. Gaussian surface about an infinite sheet of charge

D is normal to the sheet

$$D = D_z \cdot a_z$$

Apply Gauss's law

$$\begin{aligned} \rho_s \oint_S ds &= Q = \oint D \cdot ds \\ &= D_z \left[\int_{top} ds + \int_{bottom} ds \right] \end{aligned}$$

'D' has no x & y components. The surface area of the box is 'A'.

$$\rho_s \cdot A = D_z [A + A]$$

$$D_z = D = \frac{\rho_s}{2} a_z \text{ (or)}$$

$$E = \frac{D}{\epsilon_0} = \frac{\rho_s}{2\epsilon_0} a_z$$

Uniformly charged sphere :

Consider a sphere of radius 'a' with a uniform charge $\rho_o \text{ C/m}^3$. To determine D everywhere construct Gaussian surface for $r \leq a$ & $r \geq a$. Separately

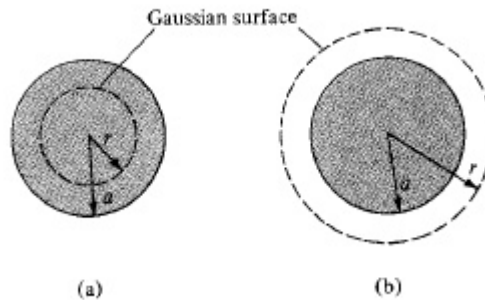


Fig27. Gaussian surface for a uniformly charged sphere a) $r \leq a$ b) $r \geq a$

For $r \leq a$. The total charge enclosed by the sphere of radius r is

$$Q_{enc} = \int_v \rho_v dv = \rho_v \int_v dv = \rho_v \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^r r^2 \sin \theta dr d\theta d\phi$$

$$Q_{enc} = \rho_v \frac{4\pi}{3} r^3 \dots\dots\dots (1)$$

and

$$\begin{aligned} \Psi &= \oint_s D \cdot ds = \oint_s D_r \cdot ds \\ &= D_r \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} r^2 \sin \theta d\theta d\phi \\ &= D_r 4\pi r^2 \dots\dots\dots (2) \end{aligned}$$

Hence $\Psi = Q_{enc}$

Equate (1) & (2)

$$D_r 4\pi r^2 = \frac{4\pi}{3} r^3 \rho_v$$

$$D_r = D = \frac{r}{3} \rho_v dr \quad 0 < r \leq a$$

For $r \geq a$ The charge enclosed by the surface is

$$Q_{enc} = \int_v \rho_v dv = \rho_v \int_v dv = \rho_v \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^r r^2 \sin \theta dr d\theta d\phi$$

$$Q_{enc} = \rho_v \frac{4\pi}{3} a^3 \quad \dots\dots\dots (3)$$

while $\Psi = \oint_s D \cdot ds = D_r \oint_s ds$

$$= D_r 4\pi r^2 \quad \dots\dots\dots (4)$$

Equate (3) & (4)

$$D_r 4\pi r^2 = \frac{4\pi}{3} a^3 \rho_v$$

$$D_r = \frac{a^3}{3r^2} \rho_v dr \quad \dots\dots\dots (5)$$

$$D = \frac{r}{3} \rho_v a_r \quad 0 \leq r \leq a$$

$$D = \frac{a^3}{3r^2} \rho_v a_r \quad r \geq a$$

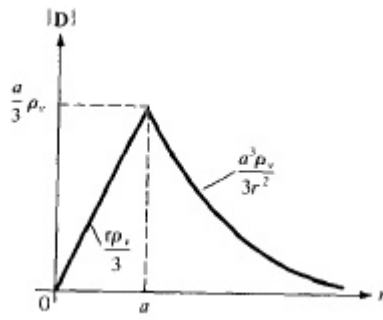


Fig28. Sketch of $|D|$ against r
for uniformly charged sphere