

#### 5.4. PARTICLE IN A FINITE POTENTIAL WELLS (Qualitative)

- Consider a particle of mass  $m$  moving with velocity  $v$  along the  $x$ -direction between  $x = 0$  and  $x = a$ .
- The walls of the box are not rigid. Hence it is represented by a potential well of finite depth.

**Step I:** Let  $E$  be the total energy of particle inside the box and  $V$  be its P.E. The potential energy which is assumed to be zero within the box and its value outside the box is finite say  $V_0$  and  $V_0 > E$ .

The variation of potential with  $x$  is shown in fig.

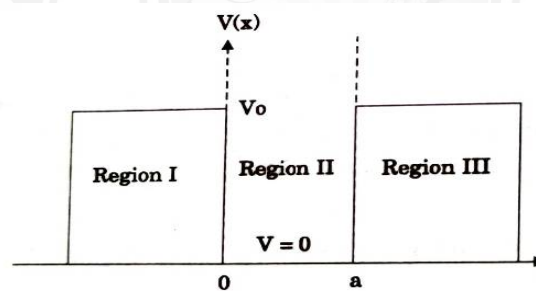


Fig. 5.11 Particle in a finite potential well with potential  $V = 0$  in the Region II enclosed between the Regions I and III with  $V = V_0$  of finite height

$$V(x) = V_0 \quad x \leq 0 \quad \text{Region I}$$

$$V(x) = 0 \quad 0 < x < a \quad \text{Region II}$$

$$\text{and } V(x) = V_0 \quad x \geq a \quad \text{Region III}$$

Classically, the particle with energy  $E < V_0$  cannot be present in regions I and III outside the box.

Consider the quantum mechanical picture of the particle in one dimension. If  $\psi$  is the wave function associated with the particle then Schrodinger's time independent equation for it is,

$$\frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2}(E - V)\psi = 0$$

**Step II:** Consider the three regions I, II, III separately and let  $\psi_I, \psi_{II}, \psi_{III}$  be the wave functions in them respectively.

We have for region I.

$$\frac{d^2\psi_I}{dx^2} + \frac{2m}{\hbar^2}(E - V_0)\psi_I = 0$$

For region II,

$$\frac{d^2\psi_{II}}{dx^2} + \frac{2mE}{\hbar^2}\psi_{II} = 0$$

( $\because V = 0$ )

and for region III,

$$\frac{d^2\psi_{III}}{dx^2} + \frac{2m(E - V_0)}{\hbar^2}\psi_{III} = 0$$

$$\frac{2mE}{\hbar^2} = k^2 \quad \text{and} \quad \frac{2m(E - V_0)}{\hbar^2} = -k'^2 \quad (\text{as } E < V_0)$$

Then the equation in the three regions is written as,

$$\begin{aligned} \frac{d^2\psi_1}{dx^2} - k^2\psi_1 &= 0 \\ \frac{d^2\psi_{1I}}{dx^2} + k^2\psi_{1I} &= 0 \\ \frac{d^2\psi_{III}}{dx^2} - k'^2\psi_{III} &= 0 \end{aligned}$$

**Step III:** The solutions of these equations are of the form.

$$\begin{aligned} \psi_1 &= Ae^{k'x} + Be^{-k'x} & \text{for } x < 0 \\ \psi_{II} &= P \cdot e^{ikx} + Q \cdot e^{-ikx} & \text{for } 0 < x < a \\ \psi_{III} &= C \cdot e^{k'x} + De^{-k'x} & \text{for } x > a \end{aligned}$$

**Step IV:** As  $x \rightarrow \pm\infty$ ,  $\psi$  should not become infinite. Hence  $B = 0$  and  $C = 0$ .

Hence the wave functions in three regions are

$$\begin{aligned}\psi_I &= Ae^{k'x} \\ \psi_{II} &= P \cdot e^{ikx} + Q \cdot e^{-ikx} \\ \psi_{III} &= D \cdot e^{-k'x}\end{aligned}$$

**Step V:** The constants  $A, P, Q$  and  $D$  can be determined by applying the boundary conditions. The wave function  $\psi$  and its derivative  $\frac{d\psi}{dx}$  should be continuous in the region where  $\psi$  is defined.

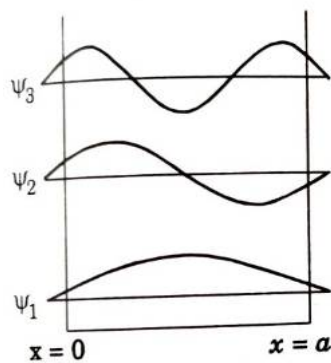
$$\psi_I(0) = \psi_{II}(0)$$

$$\left[ \frac{d\psi_I}{dx} \right]_{x=0} = \left[ \frac{d\psi_{II}}{dx} \right]_{x=0}$$

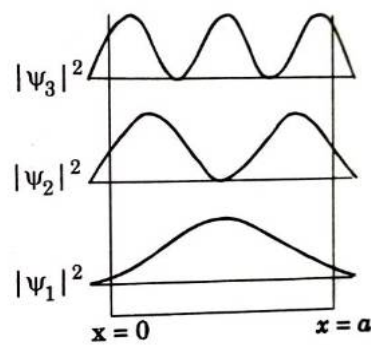
$$\psi_{II}(a) = \psi_{III}(a)$$

$$\left[ \frac{d\psi_{II}}{dx} \right]_{x=a} = \left[ \frac{d\psi_{III}}{dx} \right]_{x=a}$$

- Using these four conditions, we get four equations from which the four constants  $A, P, Q, D$  can be determined. Thus the wave functions can be known completely.
- The first three wave functions and probability densities when plotted against  $x$  are as shown in fig.



(a) Wave functions



(b) Probability densities inside non - rigid box

- The eigen functions are similar in appearance to those of infinite well except that they extend a little outside the box.
- Even though the particle energy  $E$  is less than the P.E.  $V_0$ , there is a definite probability that the particle is found outside the box.
- The particle energy is not enough to break through the walls of the box but it can penetrate the walls and leak out.
- This shows penetration of the particle into the classically forbidden region.
- The energy levels of the particle are still discrete but there are a finite number of them. Such a limit exists because, soon the particle energy becomes equal to  $V_0$ .  
For energies higher than this the particle energy is not quantised but may have any value above  $V_0$ .
- These predictions are unique in quantum mechanics and shows different behaviour from that expected in classical physics.

