

**Ampere's Circuital Law:**

Ampere's circuital law states that the line integral of the magnetic field  $\vec{H}$  (circulation of  $H$ ) around a closed path is the net current enclosed by this path. Mathematically,

$$\oint \vec{H} \cdot d\vec{l} = I_{enc} \dots\dots\dots(4.8)$$

The total current  $I_{enc}$  can be written as,

$$I_{enc} = \int_S \vec{J} \cdot d\vec{s} \dots\dots\dots(4.9)$$

By applying Stoke's theorem, we can write

$$\begin{aligned} \oint \vec{H} \cdot d\vec{l} &= \int_S \nabla \times \vec{H} \cdot d\vec{s} \\ \therefore \int_S \nabla \times \vec{H} \cdot d\vec{s} &= \int_S \vec{J} \cdot d\vec{s} \\ \therefore \nabla \times \vec{H} &= \vec{J} \dots\dots\dots(4.10) \end{aligned}$$

which is the Ampere's law in the point form.

**Applications of Ampere's law:**

We illustrate the application of Ampere's Law with some examples.

**Example 4.2:** We compute magnetic field due to an infinitely long thin current carrying conductor as shown in Fig. 4.5. Using Ampere's Law, we consider the close path to be a circle of radius  $\rho$  as shown in the Fig. 4.5.

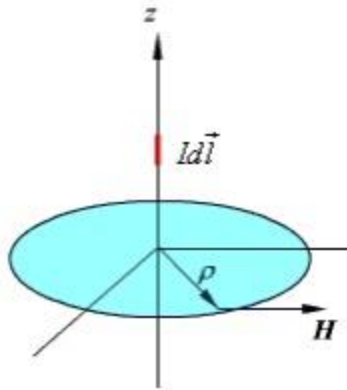
If we consider a small current element  $ld\vec{l}(= ldz\hat{a}_z)$ ,  $d\vec{H}$  is perpendicular to the plane containing both  $d\vec{l}$  and  $\vec{R}(= \rho\hat{a}_\rho)$ . Therefore only component of  $\vec{H}$  that will be present is  $H_\phi$ , i.e.,  $\vec{H} = H_\phi\hat{a}_\phi$ .

By applying Ampere's law we can write,

$$\vec{H} = \frac{I}{2\pi\rho} \hat{a}_\phi \int_0^{2\pi} H_\phi \rho d\phi = H_\phi \rho 2\pi = I \dots\dots\dots(4.11)$$

Therefore,  $\vec{H} = \frac{I}{2\pi\rho} \hat{a}_\phi$  which is same as equation (4.7)





**Fig. 4.5: Magnetic field due to an infinite thin current carrying conductor**

**Example 4.3:** We consider the cross section of an infinitely long coaxial conductor, the inner conductor carrying a current  $I$  and outer conductor carrying current  $-I$  as shown in figure 4.6.

We compute the magnetic field as a function of  $\rho$  as follows:

In the region  $0 \leq \rho \leq R_1$

$$I_{enc} = I \frac{\rho^2}{R_1^2} \dots\dots\dots(4.12)$$

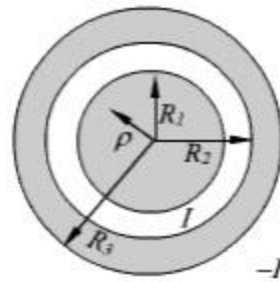
$$H_\phi = \frac{I_{enc}}{2\pi\rho} = \frac{I\rho}{2\pi a^2} \dots\dots\dots(4.13)$$

In the region  $R_1 \leq \rho \leq R_2$

$$I_{enc} = I$$

$$H_\phi = \frac{I}{2\pi\rho} \dots\dots\dots(4.14)$$





**Fig. 4.6: Coaxial conductor carrying equal and opposite currents**

In the region  $R_2 \leq \rho \leq R_3$

$$I_{enc} = I - I \frac{\rho^2 - R_2^2}{R_3^2 - R_2^2} \dots\dots\dots(4.15)$$

$$H_\phi = \frac{I}{2\pi\rho} \frac{R_3^2 - \rho^2}{R_3^2 - R_2^2} \dots\dots\dots(4.16)$$

In the region  $\rho > R_3$

$$I_{enc} = 0 \quad H_\phi = 0 \dots\dots\dots(4.17)$$